

# Gauge-Higgs Unification with Kinetic Brane Terms

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# Outline

- Introduction (EWSB and SM parameters),
- Phenomenological EW-Higgs Unification<sup>a</sup>,
- 6D Gauge-Higgs unification with  $SU(3)_W$ ,
- Brane Kinetic Terms and  $\theta_W$ <sup>b</sup>,
- Higgs mass with Brane Kinetic Terms,
- Conclusions,

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<sup>a</sup>Based on work done with A. Aranda and A. Rosado, [MPL\(2006\)](#)

<sup>b</sup>Based on work done with A. Aranda, [PLB\(2006\)](#)

# Introduction

- EWSB needed to generate SM gauge and fermion masses.
- SM Higgs doublet  $\Phi(x)$  has the potential:  
$$V(\Phi) = -\mu^2|\Phi|^2 + \lambda|\Phi|^4$$
- One unknown parameter  $\lambda$  determines  $M_H$
- Exp. and precision analysis seems to prefer a “light” Higgs boson, i.e.  $116 \leq m_h \leq 180$  GeV.
- LHC is expected to detect at least one Higgs boson, while its precise nature will be tested at LC.

# Hierarchy Problems

- **Large Hierarchy**: a severe fine tuning problem,  
Scalars get quad. rad. corr.:  $\delta M_H^2 \sim C_i \Lambda^2$   
A mechanism should protect Higgs mass  
when  $\Lambda \gg m_W$ , i.e.  $C_i \rightarrow 0$ ,
- **Little Hierarchy**  
Higgs mass should not be far from  $\Lambda_{eff}$ ,  
Global fits suggest  $\Lambda_{eff} \sim 400 \text{ GeV}$ ,  
EW precision measurements imply,  
 $\Lambda_{eff} \sim 5 - 10 \text{ TeV}$

# Models of EWSB

- Traditional solutions to EWSB problems:
  - Weak/perturbative theories, e.g. SUSY,
  - Strongly-interacting dynamics, e.g. TC
- New approaches proposed recently are based on:
  - New Dynamics: - Little/Fat Higgs,
    - String motivated.....
  - Extra Dimensions: - EWSB by O.B.C.
    - Higgsless EWSB, - Warped-AdS/CFT,
    - Gauge-Higgs unification

# SM Parameters

- Gauge parameters (**dimensionless**)
  - Gauge couplings:  $g_3, g_2, g_Y$
  - Strong phase:  $\theta_{QCD}$
- Higgs (**dimensionfull**) parameter:  $\mu^2$
- Higgs self-coupling:  $\lambda$
- Yukawa Couplings  $Y_{ij}^f$
- Could it be possible to express all SM parameters in terms of gauge parameters?

# Phenomenological Unification

# EW-Higgs

- Higgs self-coupling should be related to EW gauge couplings,
- This relation could take a generic form:  
$$\lambda = f(g_1, g_2)$$
- To simplify we impose a linear (I) or quadratic (II) relation at a high scale,
- The quadratic relation (case II) takes the form:  
$$g_i^2 = k_H \lambda,$$
- The normalization factor  $k_H$  is taken as:  
$$0.1 < k_H < 10$$

# Phenomenological Unification

# EW-Higgs

- What is the high scale? We take it where  $g_1 = g_2$ ,
- Result depends on hypercharge normalization, i.e.  $g_1^2 = k_Y g_Y^2$ ,
- We considered  $k_Y = \frac{5}{3}, \frac{3}{2}, \frac{7}{4}$ ,
- Use RGE to get  $\lambda$  at EW scale and determine the Higgs mass.
- We obtain the result:  $m_h \simeq 180$  GeV.



# XD and Gauge-Higgs Unification

- Basic idea: Identify the Higgs as a component of a higher dimensional gauge field ( $A_{\hat{\mu}}$ )

$$A_{\hat{\mu}} \rightarrow A_{\mu} = 4\text{D gauge bosons}$$

$$A_M = 4\text{D scalars}$$

- **5D Models:**  
Generally predict a very small Higgs mass
- However, in **6D Models:**  
Quartic coupling present at tree level  
Higgs mass prediction is better  
But EWSB stability must be verified

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- $F_{\hat{\mu}\hat{\nu}} = \partial_{\hat{\mu}} A_{\hat{\nu}} - \partial_{\hat{\nu}} A_{\hat{\mu}} - ig_6 [A_{\hat{\mu}}, A_{\hat{\nu}}]$

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- Spacetime (for  $N = 2$ ):  $A^\mu \rightarrow A^\mu, A^M \rightarrow -A^M, M = 5, 6$ .
- Gauge: for the generators of SU(3),  $t_a \rightarrow \Theta^{-1} t_a \Theta$ , where

$$\Theta = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix} .$$

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- $A_\mu = \sum_{a=1,2,3,8} A_\mu^{(a)} \frac{\lambda_a}{2}$
- $H_M = \sum_{a=4,5,6,7} A_M^{(a)} \frac{\lambda_a}{2} ; M = 5, 6$



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$$A_\mu = \frac{1}{2} \begin{pmatrix} A_\mu^{(3)} + \frac{1}{\sqrt{3}} A_\mu^{(8)} & A_\mu^{(1)} - i A_\mu^{(2)} & 0 \\ A_\mu^{(1)} + i A_\mu^{(2)} & -A_\mu^{(3)} + \frac{1}{\sqrt{3}} A_\mu^{(8)} & 0 \\ 0 & 0 & -\frac{2}{\sqrt{3}} A_\mu^{(8)} \end{pmatrix}$$
$$= \frac{1}{2} \begin{pmatrix} W_\mu^{(3)} + \frac{1}{\sqrt{3}} A_\mu^{(8)} & \sqrt{2} W_\mu^+ & 0 \\ \sqrt{2} W_\mu^- & -W_\mu^{(3)} + \frac{1}{\sqrt{3}} A_\mu^{(8)} & 0 \\ 0 & 0 & -\frac{2}{\sqrt{3}} A_\mu^{(8)} \end{pmatrix}$$

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$$H_M = \frac{1}{2} \begin{pmatrix} 0 & 0 & A_M^{(4)} + iA_M^{(5)} \\ 0 & 0 & A_M^{(6)} + iA_M^{(7)} \\ A_M^{(4)} - iA_M^{(5)} & A_M^{(6)} - iA_M^{(7)} & 0 \end{pmatrix}$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 & H_M^{*+} \\ 0 & 0 & H_M^0 \\ H_M^- & H_M^{0*} & 0 \end{pmatrix}$$

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$$\begin{aligned}\mathcal{L}_4^{(0)} &= -\frac{1}{4}(F_{\mu\nu}^a)^2 - \frac{1}{4}(B_{\mu\nu})^2 \\ &+ \left| \left( \partial_\mu - ig_4 W_\mu^a \frac{\tau^a}{2} - ig_4 \tan \theta_W \frac{1}{2} B_\mu \right) \mathcal{H} \right|^2 - \frac{g_4^2}{2} |\mathcal{H}|^4\end{aligned}$$

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- $\mathcal{H} = \begin{pmatrix} H_M^{*+} \\ H_M^0 \end{pmatrix}$
- $\tan \theta_W = \sqrt{3}$
- $V_{class}(\mathcal{H}) = \frac{g_4^2}{2} |\mathcal{H}|^4$



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- Arbitrary powers of  $F_{5,6}^p \rightarrow p = 1$  (quadratic)  $p = 2$  (log)
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Approximately

$$V_{quant}(\mathcal{H}) = -\mu^2|\mathcal{H}|^2 + \lambda|\mathcal{H}|^4$$

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- Suppose  $\mu^2 > 0$
- Use  $\langle |\mathcal{H}| \rangle = v/\sqrt{2}$
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$$\rightarrow \frac{m_H}{m_W} = \frac{2\sqrt{2\lambda}}{g} = 2^a$$

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<sup>a</sup>C.A. Scrucca, M. Serone, L. Silvestrini, A. Wulzer, JHEP **0402** (2004) 049; A. Wulzer, hep-th/0405168; C. Biggio, M. Quirós; hep-ph/0407348

# 6D SU(3) with Brane Kinetic Terms (BKT)

In order to fix the value of  $\tan \theta_W$  we introduce brane kinetic terms by considering the following possibilities:

- **A:**  $\mathcal{L}_{TCB} = -\frac{1}{4} c \delta(x_5) \delta(x_6) F^{\mu\nu} F_{\mu\nu}$



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- **C:**  $\mathcal{L}_{TCB} = -\frac{1}{4} \delta(x_5) \delta(x_6) \left[ c_1 \left( F_{\mu\nu}^{(a)} \right)^2 + c_2 \left( F_{\mu\nu}^{(8)} \right)^2 \right]$   
 $a = 1, 2, 3$

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C:

$$\begin{aligned}\mathcal{L}_4^{(0)} &= -\frac{1}{4}(F_{\mu\nu}^a)^2 - \frac{1}{4}(B_{\mu\nu})^2 \\ &+ \left| \left( \partial_\mu - i \frac{g_4}{\sqrt{\mathcal{Z}_1}} W_\mu^a \frac{\tau^a}{2} - i \frac{g_4}{\sqrt{\mathcal{Z}_2}} \sqrt{3} \frac{1}{2} B_\mu \right) \mathcal{H} \right|^2 \\ &- \frac{g_4^2}{2} |\mathcal{H}|^4\end{aligned}$$

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- $\mathcal{Z}_{1,2} = 1 + \frac{c_{1,2}}{\pi^2 R_5 R_6}$
- Study parameter space consistent with  $\tan \theta_W$
- Check the Higgs mass

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**Weinberg Angle:** Defining  $g = g_4 / \sqrt{\mathcal{Z}_1}$  and  $g' = \sqrt{3}g_4 / \sqrt{\mathcal{Z}_2}$

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**Higgs Mass:**  $\frac{m_H}{m_W} = 2\sqrt{\mathcal{Z}_1}$



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